

## One-atom-layer 4×4 compound in (Tl, Pb)/Si(111) system



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### ABSTRACT

An ordered 4×4-periodicity 2D compound has been found in the (Tl, Pb)/Si(111) system and its composition, structure and electronic properties have been characterized using low-energy electron diffraction, scanning tunneling microscopy observations and density-functional-theory calculations. The compound has been concluded to contain 9 Tl atoms and 12 Pb atoms per 4×4 unit cell, i.e., 0.56 ML of Tl and 0.75 ML of Pb. Structural model was proposed for the 4×4-(Tl, Pb) compound where building blocks are a hexagonal array of 12 Pb atoms, a triangular array of 6 Tl atoms and a Tl trimer. The proposed structure has a  $C_3$  symmetry and occurs in the two equivalent orientations. The electron band structure of the compound contains two metallic spin-split surface-state bands. Bearing in mind the advanced properties of the known  $\sqrt{3} \times \sqrt{3}$  2D compound in the same (Tl, Pb)/Si(111) system (i.e., combination of giant Rashba effect and superconductivity), the found 4×4-(Tl, Pb) compound is believed to be a promising object for exploration of its superconductive properties.

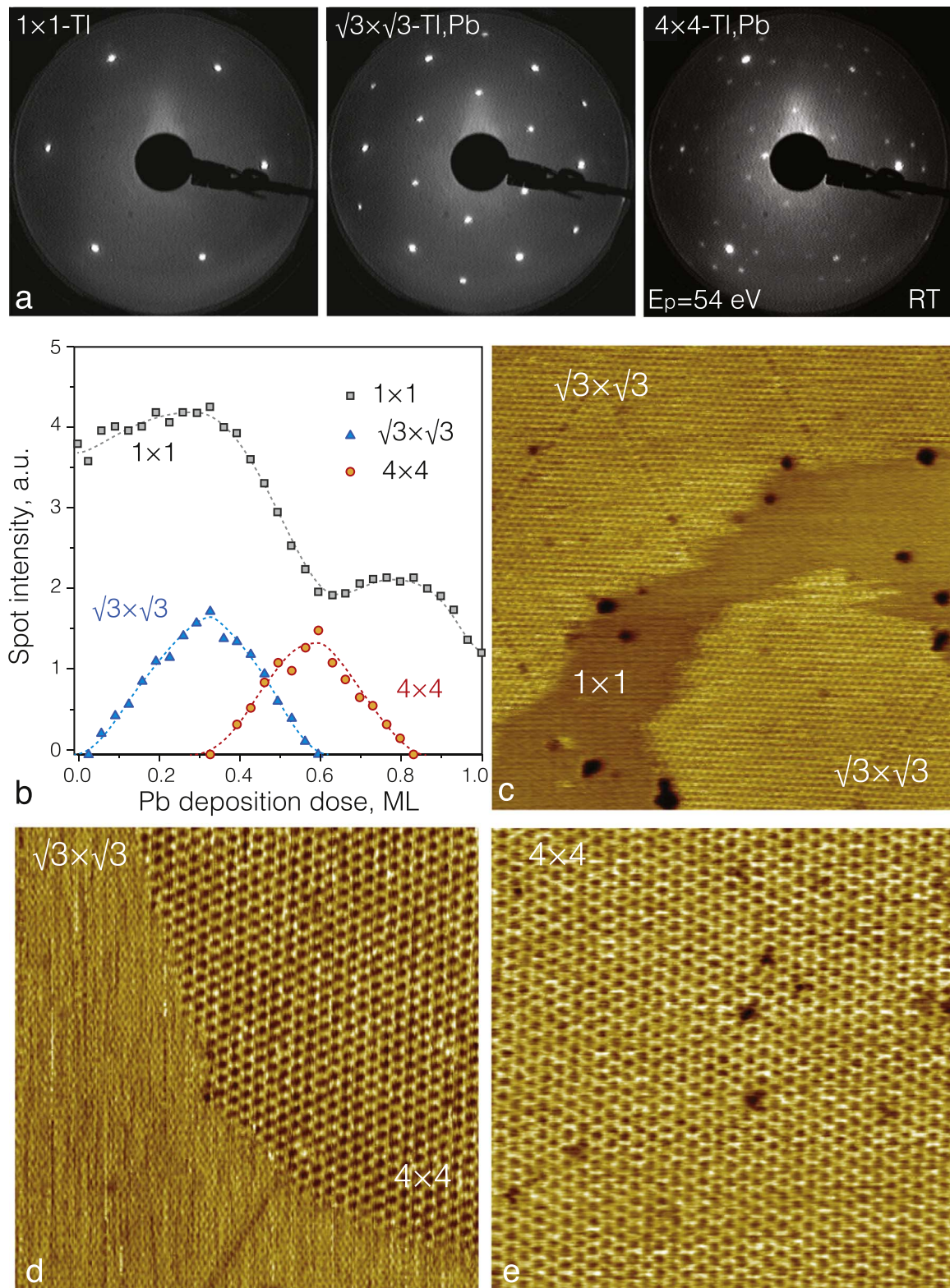
### 1. Introduction

Discovery of graphene has stimulated the current interest to atomically thin, two-dimensional (2D) materials which might exhibit unique properties not observed in the bulk materials. Metal-induced surface reconstructions on silicon and germanium (i.e., metal films of a monolayer or submonolayer on silicon or germanium) are vivid examples. Most of the studies in this field have been restricted to the reconstructions built of a single atomic species. Natural expansion of the research area is the exploration of the 2D multi-component systems. Efficiency of this approach has been proved in the very recent studies of the multi-component 2D layers on silicon [1–12] and germanium [13,14]. In particular, adding appropriate adsorbates (e.g., In, Tl, Na, or Cs) to originally structurally and electronically poor Au/Si(111)  $\sqrt{3} \times \sqrt{3}$  reconstruction converts it into a perfect 2D electron-gas system [3] with spin-split metallic surface-state bands [4,5]. A still greater promise is associated with the ordered 2D compounds or alloys which can be considered as a novel class of 2D materials [1]. For example, the (Sn, Ag)/Si(111)2×2 [1] and (Tl, Sn)/Si(111) $\sqrt{3} \times \sqrt{3}$  [11] systems show up to be metallic with steep dispersion of the surface-state bands where when crossing the Fermi level yields an electron velocity  $\sim 8 \times 10^5$  m/s that is comparable to the value of  $\sim 1 \times 10^6$  m/s reported for graphene. In addition, electronic

properties of the (Sn, Ag)/Si(111)2×2 can be tuned from metallic to semiconducting by adsorption of additional Ag onto the surface [1]. A number of 2D compounds (e.g., (Bi, Na)/Si(111) $\sqrt{3} \times \sqrt{3}$  [8,15], (Tl, Pb)/Si(111) $\sqrt{3} \times \sqrt{3}$  [8], (Au, Al)/Si(111)2×2 [10]) demonstrate a giant Rashba effect with spin splitting of the metallic bands near the Fermi level  $\Delta E_F \sim 150$ –250 meV. Among them, the (Au, Al)/Si(111)2×2 exhibits an unusual anisotropic spin texture [10]. The (Tl, Pb)/Si(111) $\sqrt{3} \times \sqrt{3}$  compound has been found to combine a giant Rashba effect and two-dimensional superconducting transport properties [8], so offers the potential to the superconducting spintronics [16].

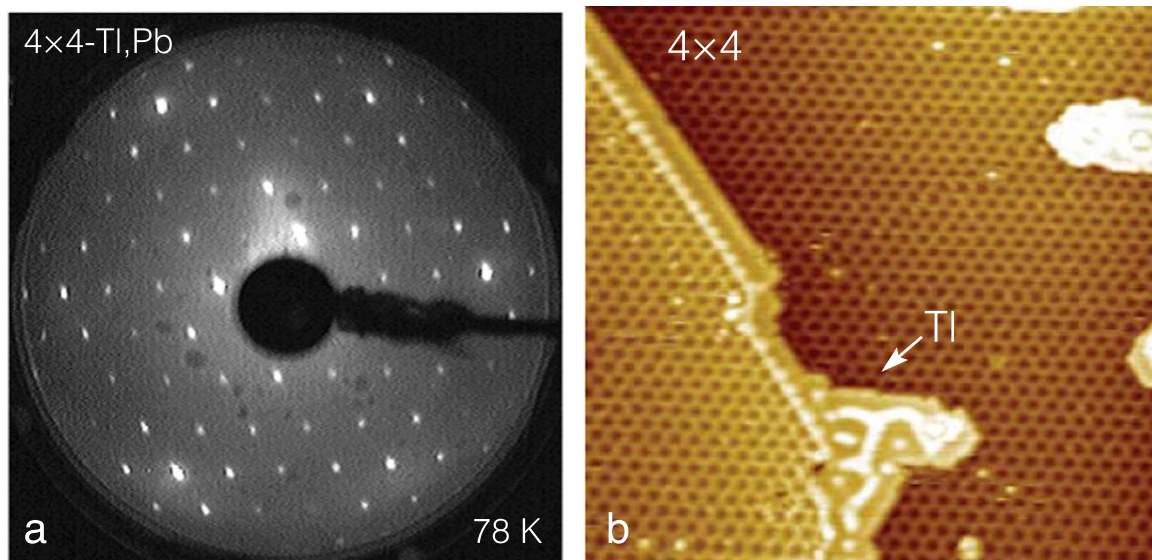
Bearing in mind these advanced properties of the (Tl, Pb)/Si(111) system, we have examined it in a more detail. In addition to the known (Tl, Pb)/Si(111)  $\sqrt{3} \times \sqrt{3}$ , one more ordered compound having 4×4 periodicity has been found. It plausibly contains 9 Tl atoms and 12 Pb atoms per 4×4 unit cell, i.e., 0.56 ML of Tl and 0.75 ML of Pb, the total metal coverage being 1.31 ML. The structural model has been derived which incorporates a hexagonal array of 12 Pb atoms, triangular array of 6 Tl atoms and a Tl trimer as the structure building blocks. The 2D compound has a  $C_3$  symmetry, hence occurs in the two possible orientations. Surmounting a barrier of 77 meV, it can twist from one orientation to another through an intermediate structure having a  $C_{3v}$  symmetry. The (Tl, Pb)/Si(111)4×4 compound has been found to be metallic with two spin-split metallic surface-state bands. The com-

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**Fig. 1.** (a) LEED patterns ( $E_p = 54$  eV) from pristine Tl/Si(111)1×1 surface and 2D (Tl, Pb)/Si(111) $\sqrt{3} \times \sqrt{3}$  and (Tl, Pb)/Si(111)4×4 compounds acquired at RT. (b) Evolution of the LEED spot intensities in the course of structural transformations induced by Pb deposition onto Tl/Si(111)1×1 surface held at RT. (c) and (d) 50×50 nm<sup>2</sup> STM images (acquired at -0.8 V, 0.6 nA and +0.5 V, 1 nA, respectively) showing the mixed  $1 \times 1/\sqrt{3} \times \sqrt{3}$  and  $\sqrt{3} \times \sqrt{3}/4 \times 4$  surfaces, respectively. (e) 50×50 nm<sup>2</sup> STM image (acquired at +1.0 V, 1.0 nA) of the (Tl, Pb)/Si(111)4×4 surface.





**Fig. 2.** (a) LEED pattern ( $E_p = 54$  eV) and (b)  $50 \times 50$  nm<sup>2</sup> STM image (+1.4 V, 0.4 nA) of the (Tl,Pb)/Si(111) $4 \times 4$  surface acquired at 78 K. The Tl double layer along the step edge is indicated by the arrow.

pound is believed to be a promising object for low-temperature transport measurements for exploration of the superconducting properties of the (Tl, Pb)/Si(111) system.

## 2. Experimental and calculation details

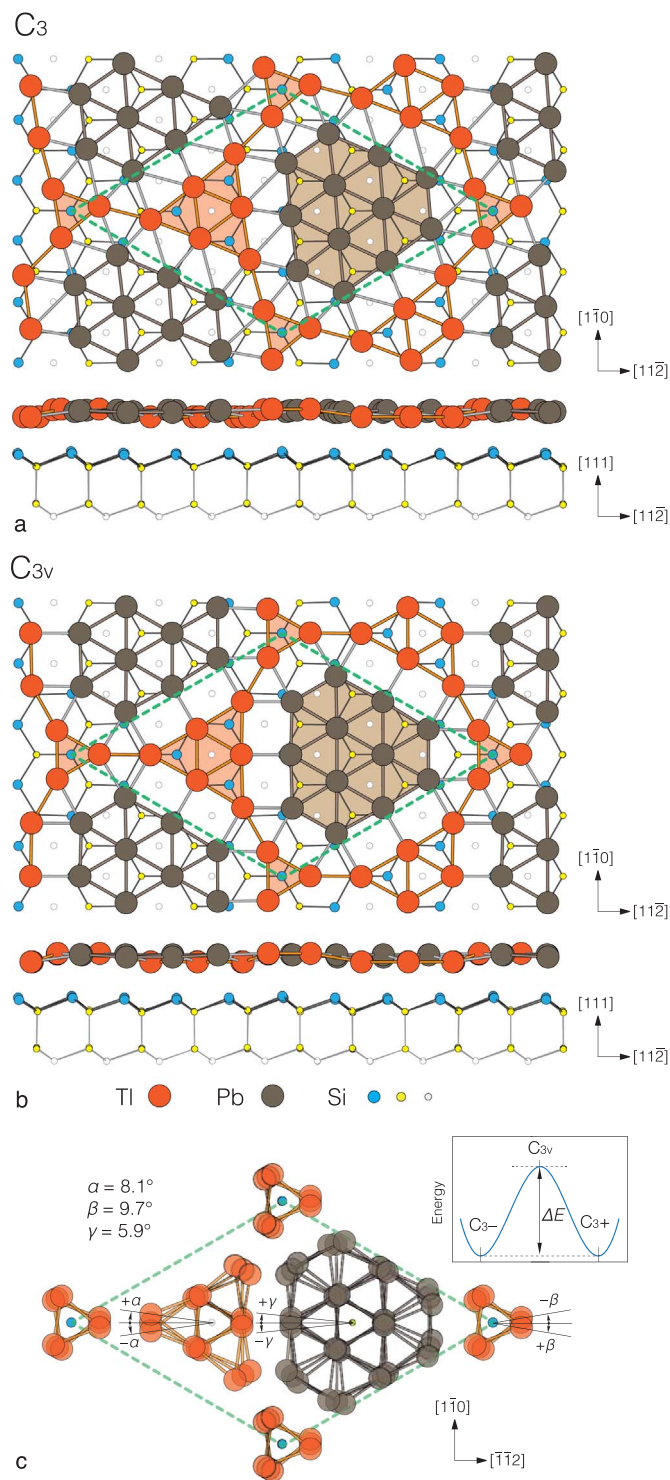
Our experiments were performed with Omicron MULTIPROBE system operated in an ultrahigh vacuum ( $\sim 2.5 \times 10^{-10}$  mbar). Atomically-clean Si(111) $7 \times 7$  surface was prepared *in situ* by flashing to 1280 °C after the sample was first outgassed at 600 °C for several hours. Tl and Pb were deposited from Ta tubes. The amounts of deposited Tl and Pb were evaluated from the deposition times using calibration based on the formation of the known reconstructions, namely Si(111) $1 \times 1$ -Tl phase containing 1.0 ML of Tl and Si(111) $\sqrt{3} \times \sqrt{3}$ -(Tl, Pb) compound incorporating 0.33 ML of Pb and 1.0 ML of Tl. Typical accuracy of such evaluations is on the order of  $\sim 20\%$ . Structure of the forming (Tl, Pb)/Si(111) surfaces was monitored using low-energy electron diffraction (LEED) and scanning tunneling microscopy (STM) observations. STM images were acquired in a constant-current mode with a mechanically cut PtIr tip after annealing in vacuum.

Our calculations were based on DFT as implemented in the Vienna *ab initio* simulation package VASP [17,18] using a planewave basis set. The projector-augmented wave approach [19] was used to describe the electron-ion interaction and the generalized gradient approximation (GGA) of Perdew, Burke, and Ernzerhof (PBE) [20] was employed as the exchange-correlation functional. The scalar relativistic effect and the spin-orbit interaction (SOI) were taken into account. To simulate the Tl-Pb reconstructions on Si(111) we used a slab consisting of four bilayers (BL) of silicon, at the PBE-optimized bulk Si lattice constant. Hydrogen atoms were used to passivate the Si dangling bonds at the bottom of the slab. The atomic positions of adsorbed atoms and atoms of Si layers within the three BLs of the slab were optimized. Silicon atoms of the deeper layers were kept fixed at the bulk crystalline positions. The kinetic cutoff energy was 400 eV, and a  $5 \times 5 \times 1$   $k$ -point mesh was used to sample the surface Brillouin zone. The geometry optimization is performed until the residual force on atoms was smaller than 10 meV/Å.

## 3. Results and discussion

Formation of the Tl-Pb compounds on the Si(111) surface is illustrated in Fig. 1 which summarizes the results of LEED and STM observations. The process was started with preparation of the high-quality Tl/Si(111) $1 \times 1$  template surface which was obtained by adsorbing 1.0 monolayer (ML) of Tl onto the Si(111) $7 \times 7$  surface held at about 300 °C [1 ML =  $7.8 \times 10^{14}$  cm<sup>-2</sup>]. The  $1 \times 1$ -Tl phase is known to contain a layer of Tl atoms occupying every  $T_4$  site on the bulk-truncated Si(111) substrate surface [21–23]. When Pb is deposited onto the Tl/Si(111) $1 \times 1$  surface held at room temperature (RT), regions of the Tl-Pb 2D compound having a  $\sqrt{3} \times \sqrt{3}$  periodicity develop (Fig. 1c). The regions merge together and cover the entire surface after deposition of  $\sim 1/3$  ML of Pb (Fig. 1b), in agreement with the previous reports [8,9]. The (Tl,Pb)/Si(111) $\sqrt{3} \times \sqrt{3}$  surface shows up in STM images as an array with a honeycomb-like inner arrangement (Fig. 1c). As determined in Ref. [8], the Tl-Pb compound layer is composed of honeycomb-chained trimers of Tl atoms with Pb atoms occupying the  $T_4$  sites in the center of each honeycomb unit. If Pb deposition is continued, the next Tl-Pb 2D compound displaying a  $4 \times 4$  periodicity starts to form (Figs. 1b and d). The  $4 \times 4$  LEED pattern demonstrates the highest intensity after deposition of slightly more than 0.6 ML of Pb (Fig. 1b). In the STM images, the (Tl, Pb)/Si(111) $4 \times 4$  surface also exhibits a honeycomb-like appearance (Fig. 1e). In the  $4 \times 4$  LEED pattern, the  $1/4$ th extra reflections are bright only in the vicinity of the main reflections (Fig. 1a). However, upon cooling the sample to 78 K, the sharp  $4 \times 4$  LEED pattern is observed where all the  $1/4$ th reflections are clearly seen (Fig. 2 a). Meanwhile, the STM appearance of the surface does not change noticeably (Fig. 2b). The only new feature is the appearance of the arrays occupied by a double Tl layer [24–26] which develop preferentially along the step edges. Occurrence of the excess Tl implies that formation of the (Tl, Pb)- $4 \times 4$  compound involves substitution of certain portion of Tl for Pb.

Following the hints that (i) Pb substitutes Tl, (ii) Pb coverage is above 0.6 ML and (iii) STM demonstrates  $C_{3v}$  symmetry of the surface, we have checked a wide set of the (Tl, Pb)/Si(111) $4 \times 4$  models of various Tl-Pb compositions using DFT calculations. The calculations prove that the structure formation energy decreases with substitution of Tl for Pb and the most optimal pathway is that when the total metal coverage remains essentially unchanged. This means that the number of the substituted Tl atoms is approximately the same as the number of



**Fig. 3.** Structure models of the Tl, Pb/Si(111)4×4 having (a)  $C_3$  and (b)  $C_{3v}$  symmetry. (c) Schematic diagram illustrating the transition between  $C_3$  structures in the two possible orientation with a  $C_{3v}$  structure as an intermediate stage.

**Table 1**  
Structural parameters of the  $C_3$  and  $C_{3v}$  models shown in Fig. 3.

Model	Tl Triangle			Tl Trimer			Pb Hexagon		
	$d_{\text{Tl-Tl}}$ , Å	$\Delta h_{\text{Tl-Si}}$ , Å	$\alpha$ , °	$d_{\text{Tl-Tl}}$ , Å	$\Delta h_{\text{Tl-Si}}$ , Å	$\beta$ , °	$d_{\text{Pb-Pb}}$ , Å	$\Delta h_{\text{Pb-Si}}$ , Å	$\gamma$ , °
$C_3$	3.26–3.39	2.39	8.1	3.10	2.86	9.7	3.23–3.48	2.75	5.9
$C_{3v}$	3.22–3.45	2.36	0	3.06	2.88	0	3.18–3.35	2.81	0

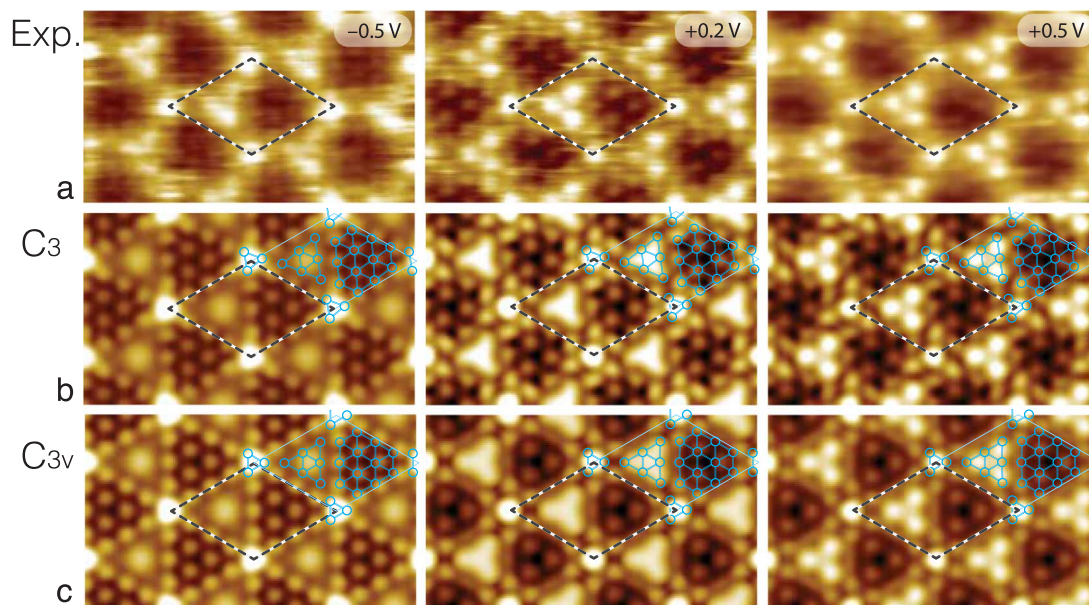
the adopted Pb atoms. Thus, it has been appeared that the most plausible Tl-Pb compounds contain 9 Tl atoms and 12 Pb atoms per 4×4 unit cell, i.e., 0.56 ML of Tl and 0.75 ML of Pb, the total metal coverage being 1.31 ML. Remind that the preceding  $\sqrt{3} \times \sqrt{3}$ -(Tl, Pb) phase contained 1.0 ML of Tl and 0.33 ML of Pb, hence the total metal coverage was 1.33 ML. One can note that the model Pb coverage of 0.75 ML exceeds the value of  $\sim 0.6$  ML where the 4×4 LEED intensity curve displays the maximum (see Fig. 1b). This can be attributed to the presence of the Tl double-layer patches like that shown in Fig. 2b. As a result, the 4×4-(Tl, Pb) compound never occupies the whole surface area and its Pb contents is greater than the averaged density of Pb atoms at the surface.

Among the  $C_{3v}$  structure models with the optimal composition, the one shown in Fig. 3b has the lowest formation energy. However, if the system is allowed to relax in the absence of mirror-symmetry limitations it evolves to the one having  $C_3$  symmetry as shown in Fig. 3a. Formation energy of the  $C_3$  model is lower than that of the  $C_{3v}$  model by  $\sim 77$  meV per 4×4 unit cell when PBE approximation is used. For LDA approximation this value is somewhat lower and is  $\sim 50$  meV. Both models have a very similar atomic arrangement being composed of the same building blocks. In the models, 12 Pb atoms form a hexagonal array in which Pb atoms occupy positions close to the on-top ( $T_1$ ) sites. The Pb-atom array has a shape of the hexagon which can be visualized also as a truncated triangle. Nine Tl atoms are arranged into the two structural elements. The first element contains 6 Tl atoms forming an array of triangular shape where Tl atoms occupy positions close to the  $T_4$  sites. The other 3 Tl atoms form a trimer centered in the  $T_1$  site while Tl atoms are located close to the  $H_3$  sites. Table 1 shows structural parameters of the two models. Lacking of the mirror symmetry in the  $C_3$ -model is due to the twist of structural elements, namely, the Pb-atom array is twisted by  $5.9^\circ$ , the Tl six-atom triangle by  $8.1^\circ$  and the Tl trimer by  $9.7^\circ$  with respect to the orientations of their counterparts in the  $C_{3v}$ -model. Fig. 3c illustrates occurrence of the  $C_3$ -model structure in the two equivalent orientations with twisting of structural elements either clockwise or counter-clockwise. One can also see in Fig. 3c that transition of the  $C_3$ -structure from one orientation to another occurs through the  $C_{3v}$ -structure as an intermediate stage. The transition implies surmounting a barrier. Remind that the maximal barrier value obtained within PBE approximation is  $\sim 77$  meV. Nevertheless, this is a relatively small value and one could expect that the structure is in a continuous twisting between two orientations even at relatively low temperatures.

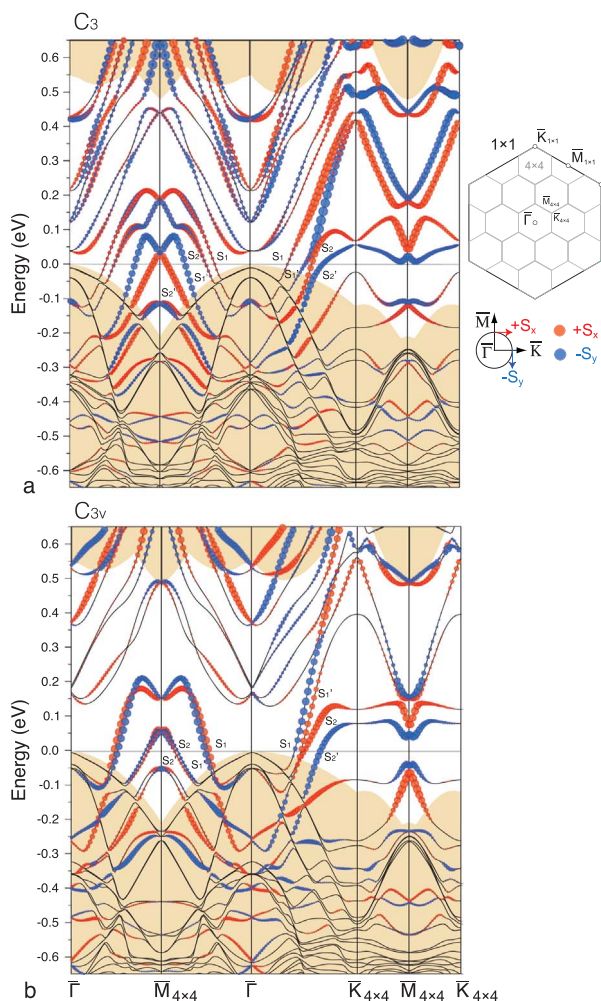
As a test for the validity of the proposed models, we have compared simulated STM images obtained for the models with the experimental high-resolution STM images where the fine structural features can be seen. The results of the comparison are summarized in Fig. 4. One can see that the simulated STM images for the  $C_3$ - and  $C_{3v}$ -models (Figs. 4 b and c) are almost identical as a natural sequence of their structural similarity. In turn, both models demonstrate a close resemblance with the experimental STM images (Fig. 4a) that sounds supportive for the correct choice of the models.

Having plausible structure models, we have calculated their electron band structures shown in Fig. 5. Fig. 5a shows the calculated band structure for the  $C_3$ -model and Fig. 5b shows that for the  $C_{3v}$ -model. The band structures are represented by the dispersion curves along the  $\bar{\Gamma}$ - $\bar{M}$ - $\bar{\Gamma}$  and  $\bar{\Gamma}$ - $\bar{K}$ - $\bar{M}$ - $\bar{K}$  directions of the 4×4 surface Brillouin zone





**Fig. 4.** Comparison of (a) experimental STM images acquired at  $-0.5$  V,  $+0.2$  V and  $+0.5$  V biases with the corresponding simulated STM images for (b)  $C_3$ - and (c)  $C_{3v}$ -models of the (Tl, Pb)/Si(111) $4\times 4$  surface. The  $4\times 4$  unit cells are outlined. The main structural elements of the models (Pb-atom array, Tl-atom triangular array and Tl trimers) are superposed onto the simulated STM images to illustrate the origin of the STM features.



**Fig. 5.** Electron band structures calculated for (a)  $C_3$ - and (b)  $C_{3v}$ -models of the (Tl, Pb)/Si(111) $4\times 4$  phase.

(SBZ) as indicated in the schematics in Fig. 5. One can see that both structures contains two spin-split metallic surface-state bands, denoted  $S_1$  ( $S_1'$ ) and  $S_2$  ( $S_2'$ ). Remarkably, spin splitting in the case of the  $C_3$ -model structure is noticeably greater than that for the  $C_{3v}$  structure (see Table 2). It is worth noting also that despite the difference in the periodicity the calculated band structure of the  $4\times 4$ -(Tl, Pb) compounds is qualitatively akin those calculated and experimentally determined with ARPES for the  $\sqrt{3} \times \sqrt{3}$ -(Tl, Pb) compound [8].

#### 4. Conclusions

In conclusion, we have extended the list of the 2D compounds in the (Tl, Pb)/Si(111) system by finding the Tl-Pb compounds having a  $4\times 4$  periodicity. Composition, atomic structure and electronic properties of the  $4\times 4$ -(Tl, Pb) compound has been characterized using combination of LEED and STM observations and DFT calculations. The compound has been concluded to contain 9 Tl atoms and 12 Pb atoms per  $4\times 4$  unit cell, i.e., 0.56 ML of Tl and 0.75 ML of Pb, the total metal coverage being 1.31 ML. Its atomic model has been proposed which is built of three basic blocks, including a hexagonal array of 12 Tl atoms, triangular array of 6 Pb atoms and a Pb trimer. The  $4\times 4$ -(Tl, Pb) compound occurs in the two equivalent orientations and is plausibly twisting between them even at relatively low temperatures, taking into account a low barrier of  $\sim 77$  meV. Compound has been found to be metallic with two spin-split metallic surface-state bands.

We believe that the (Tl, Pb)/Si(111) compound is a promising object for low-temperature transport measurements for exploration of the superconducting properties of the (Tl, Pb)/Si(111) system. In particular, it allows to address the effect of the Tl-Pb composition and structure on the properties of the 2D-compound superconductor. Compared to the  $\sqrt{3} \times \sqrt{3}$ -(Tl, Pb), the  $4\times 4$ -(Tl, Pb) compound is more Pb-rich. However, it is difficult to suggest in advance, if this would enhance superconducting properties or not. Remind that the critical temperature for the  $\sqrt{3} \times \sqrt{3}$ -(Tl, Pb) compound is 2.3 K [9], for the Pb films thicker than 8 ML it approaches a bulk value of  $\sim 7$  K [27], for the one-atom-layer Pb film on Si(111) the transport measurements yield 1.1 K [28], while the STS-based measurements yield 1.52 K and 1.83 K, depending on the structure and density of Pb-atom layer [29].

**Table 2**Calculated values of the spin split parameters of  $S_1$  and  $S_2$  zones in the  $\bar{\Gamma} - \bar{M}$  and  $\bar{\Gamma} - \bar{K}$  directions in the vicinity of Fermi level.

Model	$\bar{\Gamma} - \bar{M}$				$\bar{\Gamma} - \bar{K}$			
	$S_1$		$S_2$		$S_1$		$S_2$	
	$\Delta k, \text{\AA}^{-1}$	$\Delta E, \text{meV}$	$\Delta k, \text{\AA}^{-1}$	$\Delta E, \text{meV}$	$\Delta k, \text{\AA}^{-1}$	$\Delta E, \text{meV}$	$\Delta k, \text{\AA}^{-1}$	$\Delta E, \text{meV}$
$C_3$	0.023	55	0.047	105	0.045	77	0.027	52
$C_{3v}$	0.009	34	0.008	17	0.017	48	0.042	128

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